Cr and Na atoms, can be seen in Fig. 1. The two CrO_4 tetrahedra, inside each unit cell, are disposed with their apices in opposite directions along the *c* axis (Figs. 1 and 2). This kind of packing is the same as that of the isomorphous compound $K_3Na(SO_4)_2$ (Okada & Ossaka, 1980).

The Na—O distances in the NaO₆ octahedron are symmetry restricted to a single value of 2.364 Å. The mean distances in the K(1)O₁₀ and K(2)O₁₂ polyhedra are 2.969 and 3.155 Å respectively.

The Cr—O distances range from 1.613 to 1.6459 Å with a mean value of 1.639 Å. The values of the two possible bond angles O-Cr-O, in the CrO₄ group $(110.31 \text{ and } 108.63^\circ)$, indicate a very small distortion with respect to the ideal tetrahedral configuration. The rigidity of this group was revealed by means of a librational thermal-motion analysis based on the anisotropic atomic displacement parameters of the Cr. O(1) and O(2) atoms. The results show that the thermal motion of the CrO₄ groups can be completely described by the T and L tensors (Schomaker & Trueblood, 1968). The residual deformation deduced from the comparison between the observed and calculated atomic displacement parameters is practically zero for this tetrahedron. The corrected bond lengths involving Cr-O(1) and Cr-O(2) atom pairs (see Table 2) exhibit, therefore, closer values (1.646 and 1.661 Å respectively).

The elastic stiffness tensor component C_{44} shows a usual decrease with increasing temperature from about 250 K up to room temperature. Below 250 K, in the vicinity of the phase transition, the C_{44} component strongly decreases and acoustic phonon softening is quite evident (Gómez-Cuevas & Breczewski, 1990). The strong damping of ultrasound waves by the ferroelastic domain walls prevents measurement below the phase transition.

In the present case, the softening of C_{44} is necessary only for the appearance of 2/m symmetry in the ferroelastic phase (Toledano, Fejer & Auld, 1983). In

fact, such a symmetry change, as probed by ultrasonic measurements, would also require similar behaviour in $C_{11}-C_{12}$ (equal to $2 \times C_{66}$ for this point group) at the transition temperature. Unfortunately, the morphology of the samples hindered the study of $C_{11}-C_{12}$ as a function of temperature.

Nevertheless, previous results published by Krajewski, Mroz, Piskunowicz & Breczewski (1990) concerning the domain structure of the ferroelectric phase, together with the X-ray measurements presented above and the behaviour of the elastic tensor component C_{44} confirm, as expected from Landau theory (Toledano, Fejer & Auld, 1983; Sapriel, 1975), the change of point symmetry in the phase transition from 3m to 2/m.

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Structure of Li₄Ge₅O₁₂ – a New Compound in the Li₂O–GeO₂ System

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Abstract. Tetralithium pentagermanate Li₄Ge₅O₁₂, $M_r = 582.7$, triclinic, $P\overline{1}$, a = 5.120 (1), b = 9.143 (2), c = 9.586 (2) Å, $\alpha = 72.95$ (3), $\beta = 77.74$ (3), $\gamma =$ 78.81 (3)°, V = 415.02 (15) Å³, Z = 2, $D_m = 4.63$, $D_x = 4.663$ Mg m⁻³, λ (Mo K α) = 0.71069 Å, $\mu = 17.74$ mm⁻¹, F(000) = 536, T = 295 (2) K, R =© 1990 International Union of Crystallography

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0.025 for 1334 independent observed reflections. The structure consists of quasi-close-packed planes of O atoms held together by Ge and Li cations. Two of the five Ge in the asymmetric unit are tetrahedrally coordinated with $\langle Ge-O \rangle = 1.748$ (21) Å and the other three are octahedrally coordinated with $\langle Ge-O \rangle = 1.901$ (59) Å. Crystals up to 1.5 cm in diameter and 4.0 cm in length were pulled from a melt of composition $3Li_2O.8GeO_2$ at 1208 (5) K.

Introduction. Recently, optical studies of a Ti^{4+} -activated lithium germanium oxide compound were reported (Loiacono, Shone, Mizell, Powell, Quarles & Elouadi, 1986). The composition of the host crystal was thought to be 'Li₆Ge₈O₁₉'. This was based on the phase diagram of Murthy & Ip (1964) who reported a congruently melting phase of this composition. Subsequent studies show that under the growth conditions described below, crystals of composition Li₄Ge₅O₁₂ grow from a melt of composition Li₆Ge₈O₁₉. Li₄Ge₅O₁₂ is a new crystalline phase in the Li₂O-GeO₂ system.

The Li₂O-GeO₂ system has been reported by Murthy & Ip (1964) to contain five congruently melting compounds. These are listed in Table 1 along with other compounds that have been isolated as single crystals. Of the five congruently melting phases, four have been isolated and studied as single crystals. These are: Li₄GeO₄, Li₆Ge₂O₇, Li₂GeO₃ and $Li_2Ge_7O_{15}$. Some doubt has been expressed about the existence of 'Li₆Ge₈O₁₉' (Wittman & Modern, 1965; Ikeda & Sato, 1984). This doubt is supported by our observation that 24 of the first 25 X-ray powderdiffraction lines listed by Murthy & Ip (1964) as characteristic of 'Li₆Ge₈O₁₉' correspond quite closely to lines from $Li_4Ge_5O_{12}$ (given in Table 2), and/or Li2GeO3 (JCPDS No. 34-6590), and/or Li2Ge4O9 (JCPDS No. 37-1363), and/or Li₂Ge₇O₁₅ (JCPDS No. 23-1180) (Powder Diffraction File, JCPDS International Centre for Diffraction Data, Swarthmore, Pennsylvania, USA).

Experimental.

Crystal growth

The initial intent of these experiments was to grow crystals of composition 'Li₆Ge₈O₁₉'. Such crystals would contain 90.32 wt % GeO₂. With that goal in mind the procedure described below was used and large single crystals were obtained. These crystals were later shown by an X-ray single-crystal structure determination to be $Li_4Ge_5O_{12}$ which contains 89.75 wt % GeO₂.

Crystal growth of $Li_4Ge_5O_{12}$ proceeded as follows. AESAR, 99.99% pure Li_2CO_3 and hexagonal GeO_2 were the starting materials. These were mixed in the molar ratio 3:8, ball milled using an alumina conTable 1. Li-Ge-O phases reported to be congruently melting by Murthy & Ip (1964) and phases for which single-crystal structure determinations have been made

It is doubtful that the phase labeled Li₆Ge₈O₁₉ exists.

| Phase | Ratio of Li ₂ O:GeO ₂ | Wt % GeO ₂ | Congruently melting | Space group | Structure reference |
|---|--|--------------------------|---------------------|-------------------|--------------------------------|
| Li₄GeO₄ | 2:1 | <u>63</u> .64 | Yes | Bmmb | Vollenkle & Wittman (1969) |
| Li ₆ Ge ₂ O ₇ | 3:2 | 70.00 | Yes | $P2_1/n$ | Vollenkle (1980) |
| Li2GeO3 | 1:1 | 77.78 | Yes | Cmc2 ₁ | Vollenkle & Wittman (1968a) |
| Li2Ge2O3 | 1:2 | 87.50 | No | Сс | Vollenkle & Wittman (1968b) |
| Li ₄ Ge ₅ O ₁₂ | 2:5 | 89.75 | No | ΡĪ | This work |
| (Li6Ge8O19 |) 3:8 | 90.32 | (Yes) | | |
| Li2Ge4O9 | 1:4 | 93-33 | No | Pcca | Wittman & Modern (1965) |
| Li4Ge9O20 | 2:9 | 94.03 | No | <i>C</i> 2 | Vollenkle & Witman (1971) |
| Li2Ge2O15 | 1:7 | 96.08 | Yes | Pbcn | Vollenkle & Wittman (1970) |

Table 2. X-ray powder-diffraction data for Li₄Ge₅O₁₂

Pearson symbol *TP*2. Philips 17 cm vertical diffractometer used with Cu K α radiation, $\lambda = 1.54060$ Å, adjustable divergence slit maintaining a 12.5 mm sample illumination length, 0.2 mm rec. slit, 1 Soller slit, graphite diffracted-beam monochromator, scintillation detector and instrumental profile breadth of 0.10°. Data were collected at 295 K from a thin sample with particle size < 325 mesh. Intensities were measured as peak heights from 0.02° step scans and a Si external 2 θ standard was used.

| $2\theta_{exp}(^{\circ})$ | I/I _o | d | hkl |
|---------------------------|------------------|-------|-------------|
| 9.79 | 22 | 9.03 | 001 |
| 10.21 | 2 | 8.65 | 010 |
| 15.91 | 7 | 5.56 | 01 |
| 17.88 | 5 | 4.96 | 100 |
| 18.92 | 82 | 4.69 | 101, 111 |
| 19-33 | 26 | 4.59 | 110 |
| 19.63 | 100 | 4.521 | 012, 002 |
| 20.27 | 12 | 4.379 | 021 |
| 20.50 | 5 | 4.330 | 020 |
| 21.84 | 21 | 4.065 | 10 1 |
| 23.87 | 12 | 3.724 | 111, 121 |
| 24.46 | 3 | 3.637 | 022, 012 |
| 25.00 | 7 | 3.559 | 021 |
| 25.33 | 18 | 3.514 | 120 |
| 28·6 7 | 85 | 3.111 | 013 |
| 29-21 | 8 | 3.055 | 120 |
| 29·6 1 | 59 | 3.015 | 003 |
| 30.06 | 12 | 2.970 | 121 |
| 30.15 | 10 | 2.962 | 12 <u>1</u> |
| 31.36 | 6 | 2.850 | 023 |
| 31.67 | 8 | 2.823 | 121, 131 |
| 32.09 | 40 | 2.787 | 103, 032 |
| 32.86 | 7 | 2.723 | 132 |
| 33-64 | 1 | 2.662 | 130 |
| 33.99 | 4 | 2.635 | 013 122 |

tainer and balls, then calcined in air. The calcination sequence was 873 K for 17 h then 1073 K for 17 h and finally 1173 K for 17 h. The material was ball milled between firings and pure platinum crucibles were used for all calcinations and for subsequent crystal growth. Weight-loss measurements showed the amount of CO_2 evolved after the 1173 K firing corresponded to the reaction

$$3Li_2CO_3 + 8GeO_2 \rightarrow Li_6Ge_8O_{19}' + 3CO_2$$

It was from this material that large single crystals were grown using the Czochralski method and a resistance furnace. The growth temperature was maintained at 1208 (5) K and the temperature gradient of the melt/air interface was controlled to 10 K cm⁻¹. Crystals were pulled in air at rates from 0·1 to 5 mm h⁻¹ with rotation rates of 5 to 50 revmin⁻¹. Best results were obtained at 1 mm h⁻¹ and 25 rev min⁻¹ from melt volumes of 100 cm³ and yielded single-crystal boules 1·5 cm in diameter and 4·0 cm in length. Fig. 1 shows one such crystal. Grown crystals were cooled to room temperature at 50 K h⁻¹.

Concentrations of Ti⁴⁺ from 0.01 to 0.3 wt % could be incorporated in the Li₄Ge₅O₁₂ without a change in the appearance of the crystals when examined with an optical microscope. These crystals exhibited broad-band luminescence from 388 to 524 nm when excited with 266 nm radiation (Loiacono *et al.*, 1986).

Crystal structure determination

The composition and crystal structure of $\text{Li}_4\text{Ge}_5\text{O}_{12}$ were determined from room-temperature X-ray diffraction measurements of a 0.35 mm spherical crystal using a Nicolet single-crystal diffractometer. Variable speed ω scans were made of reflections in the 2θ range $3-50^\circ$ ($|h| \le 6$, $|k| \le 10$, $|l| \le 11$) using graphite-monochromated Mo Ka radiation. Two standard reflections were measured after every 48 reflections. These showed a random variation in intensity of $\pm 1.7\%$ with time. Twenty five reflections in the 2θ range $9-28^\circ$ were used to establish the lattice parameters. The measured



Fig. 1. Single crystal of $Li_4Ge_5O_{12}$ grown at 1208 (5) K by the Czochralski technique using a resistance-heated furnace. The crystal is shown mounted for fabrication. The growth direction, parallel to the ruler, is close to the crystal *a* axis.

density was determined by weighing and measuring the dimensions of a rectangular, single-crystal parallelepiped.

The crystal structure was solved by direct methods using the random tangent refinement procedures, RANT in the SHELXTL-Plus software package (Sheldrick, 1987). 2918 reflections were measured, 1459 were independent ($R_{int} = 2.12\%$) and 1334 were observed [*i.e.* $F > 4.0\sigma(F)$]. A full-matrix leastsquares refinement was carried out minimizing the quantity $\sum w(F_o - F_c)^2$ and employing the weighting scheme $w^{-1} = \sigma^2(F) + 0.001F^2$ where $\sigma(F)$ is based on counting statistics. Refinements of F_c were made by varying 191 structural parameters including anisotropic Ge temperature factors, an overall scale factor and an extinction parameter, x, which refined to 0.0081 (4) where $F^* = F[1 + 0.002xF^2/\sin(2\theta)]^{-1/4}$. No absorption correction was applied. Results of the refinement for observed F's are: R = 2.51%, wR =3.69%; and for all F's: R = 2.68%, wR = 4.72%. The largest mean $\Delta/\sigma = 0.007$, the largest difference peak is $0.77 \text{ e} \text{ Å}^{-3}$ and the largest difference hole is 0.74 e Å-3. Atomic scattering factors from International Tables for X-ray Crystallography (1974, Vol. IV) were used with anomalous-dispersion corrections for Ge and O.

Discussion. Final atomic positions and equivalent isotropic thermal parameters are given in Tables 3 and 4[†] and indicate that the structure is centrosymmetric, PI. Consistent with this symmetry, no optical second harmonic signal could be measured using an apparatus sensitive to a signal equal to $10^{-3} \times \alpha$ quartz (Dougherty & Kurtz, 1976). The structure consists of quasi-close-packed planes of O anions held together by Li and Ge cations. These oxygen planes lie very close to the ab* plane of the crystal and are stacked in an f.c.c. sequence. The planar arrangement of O atoms can clearly be seen in Fig. 2 which shows the structure as viewed along the *a* axis. The structure can also be considered in terms of Ge-O coordination polyhedra which may be linked at corners and/or along edges. The packing of these polyhedra is illustrated in Fig. 3 which is also a view along the *a* axis.

Four of the ten Ge per unit cell are in tetrahedral coordination and the other six in octahedral coordination. Taken together the tetrahedra have $\langle Ge-O \rangle = 1.748$ (21) Å, 1.723 < Ge-O < 1.778 Å and $103.8 < O-Ge-O < 113.8^{\circ}$. Each tetrahedron is corner linked to one other tetrahedron in the unit cell.

[†] Lists of structure factors and anisotropic thermal parameters have been deposited with the British Library Document Supply Centre as Supplementary Publication No. SUP 53094 (6 pp.). Copies may be obtained through The Technical Editor, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

2.153 (11) 2.074 (11)

2.315 (15)

2.296 (11)

100·1 (4) 101·0 (4)

179.8 (5)

175.6 (5)

177.4 (7)

97.1 (4)

99.6 (4)

80.3 (4)

80·1 (3)

79.0 (4)

80.1 (5)

95.2 (5)

80.0 (4)

99·0 (4)

99.4 (4)

79.0 (4) 103-6 (5)

Li(4) - O(3)Li(4) - O(5)

Li(4) - O(9A)Li(4) - O(12B)

O(3)—Li(1)—O(4*A*) O(3)—Li(1)—O(7*B*) O(3)—Li(1)—O(8*B*)

O(6)-Li(1)-O(7) O(6)-Li(1)-O(4A)

O(6)-Li(1)-O(7B)O(6)-Li(1)-O(7B)O(6)-Li(1)-O(8B)O(7)-Li(1)-O(4A)

O(7)—Li(1)—O(7B)O(7)—Li(1)—O(8B)

O(4A)-Li(1)-O(7B) O(4A)-Li(1)-O(8B)

O(7B)-Li(1)-O(8B)

O(1)-Li(2)-O(2) O(1)-Li(2)-O(1B) O(1)-Li(2)-O(1B) O(1)-Li(2)-O(3A)

O(1) - Li(2) - O(4B)

Table 3. Fractional atomic coordinates and equivalent isotropic displacement coefficients

Table 4 (cont.)

| | 4 | 4 | | 55 | |
|--|---------------------|----------------|--------------------------------|-----------------|---------------------------|
| U_{eq} is defined as one third of the trace of the orthogonalized U_{ii} tensor. | | | | | |
| | r | v | | 7 | U_{-} (Å ²) |
| $C_{\alpha}(1)$ | 0.1900 (1) | 0.2520 / | n i | 0.0252 (1) | |
| $C_{\alpha}(1)$ | 0.1007 (1) | 0.2339 (| | (1) | 0.013 (1) |
| Ge(2) | 0.4002 (1) | 0.0996 (| | 0.6452(1) | 0.013(1) |
| Ge(3) | 0.2614(1) | 0.0060 (| | 1.2692(1) | 0.013(1) |
| Ce(4) | 0.1170 (1) | 0.5042 (| 1) | 1.2003(1) | 0.013(1) |
| | 0.2674 (9) | 0.5445 (| | 0.0311(1) | 0.013(1) |
| O(1) | 0.2074 (8) | 0.3003 (| 4) | 0.4043 (4) | 0.017(1) |
| O(2) | 0.3191 (8) | 0.2031 (| 4) | 0.3069(4) | 0.017(1) |
| | 0.0542 (0) | 0.3791 (| 4) | 0.7304(4) | 0.015(1) |
| 0(4) | -0.0511(8) | 0.7305 (| 4) | 0.6004(4) | 0.015(1) |
| 0(5) | -0.0311 (8) | 0.7393 (| 4) | 0.0334 (4) | 0.010(1) |
| 0(0) | 0.0120 (8) | 0.2704 (| (4) | 0.8102(4) | 0.017(1) |
| O(7) | 0.1475 (0) | 0.4232 (| (4) (1) | 1.0120(4) | 0.015(1) |
| | 0.1473 (8) | 0.7322 (| 4) | 0.9309(4) | 0.015(1) |
| 0(9) | 0.4372 (7) | -0.0018 (| 4) | 0.0013(4) | 0.015(1) |
| | 0.1025 (7) | 0.1015 (| (4) | 0.8004(4) | 0.010(1) |
| | 0.0406 (7) | 0.1008 (| (4) | 1.0910(4) | 0.017(1) |
| U(12) | - 0.0920 (8) | 0.0039 (| (4) | 0.0200(4) | 0.010(1) |
| Li(1) | 0.1428 (15 | 0.4071 | (11) | 0.8893(10) | 0.024 (4) |
| L(2) | 0.1428 (22 | .) 0.40/1 (| (12) | 0.4036(10) | 0.027 (4) |
| $L_1(3)$ | -0.2748 (19 | | (10) | 0.9763(10) | 0.022(3) |
| 21(1) | 00107 (22 | ,, 0,,,00, | (12) | 0 0007 (12) | 0000(1) |
| Table 4. Selected bond lengths (Å) and angles (°) for | | | | | |
| | | | C ₅ O ₁₂ | | |
| Ge(1)- | ·O(4) 2·0 | 00 (4) | Li(1) | D(3) | 2.070 (10) |
| Ge(1)- | -O(6) 1·8 | 29 (4) | Li(1)— | O(4A) | 2.135 (9) |
| Ge(1)- | -O(7) 1·9 | 65 <u>(</u> 5) | Li(1)(| ວໄຄ໌ | 2.062 (13) |
| Ge(1)- | O(8A) = 1.8 | 37 (4) | Li(1)— | O(7B) | 2.141 (13) |
| Ge(1) | -O(10) 1·8 | 65 (5) | Li(1)(| D(7) | 2.213 (9) |
| Ge(1)- | -O(11) 1.8 | 61 (4) | Li(1) | O(8B) | 2.069 (10) |
| Ge(2) | -O(2) 1.7 | 23 (4) | Li(2)(| D(1) | 2.089 (14) |
| Ge(2)- | -O(9) 1.7 | 77 (4) | Li(2) | O(1 <i>B</i>) | 2.150 (11) |
| Ge(2) | -O(10) 1.7 | 32 (4) | Li(2) | D(2) | 1.995 (10) |
| Ge(2)- | -O(12) 1·7 | 54 (5) | Li(2) | O(3A) | 2.239 (11) |
| Ge(3)- | -O(1) 1.8 | 91 (4) | Li(2) | O(4 <i>B</i>) | 2.197 (10) |
| Ge(3)- | -O(1A) = 1.8 | 48 (4) | Li(2) | O(5B) | 2.051 (14) |
| Ge(3)- | -O(2) 1·9 | 60 (5) | Li(3)— | O(6A) | 2.199 (9) |
| Ge(3) | -O(3) 1-9 | 11 (5) | Li(3)(| D(8A) | 2·139 (12) |
| Ge(3)- | -O(4) 1·9 | 78 (4) | Li(3)— | D(10) | 1.997 (9) |
| Ge(3)- | -O(6) 1.8 | 61 (4) | Li(3)(| D(10A) | 2.162 (9) |
| Ge(4)- | -O(5A) 1.7 | 32 (4) | Li(3)- | D(11A) | 1.988 (10) |
| Ge(4)- | -O(9 <i>B</i>) 1.7 | 50 (4) | Li(3)— | D(11 <i>B</i>) | 2.154 (12) |
| Ge(4)- | -O(11) 1.7 | 41 (4) | Li(4)(| D(1) | 2.232 (14) |
| Ge(4) | O(12A) 1.7 | 78 (4) | Li(4)(| O(2A) | 2.036 (11) |
| | and and the second | a a 11.0 | | * | |

1.821 (4)

1.945 (5)

1.893 (4)

1.946 (4)

1.979 (4)

1.836 (5)

83·1 (4) 81·7 (2)

95.2 (2)

173·7 (2) 89·6 (2)

91.6 (2)

94·6 (2) 92·5 (2) 171·6 (2)

172.7 (2)

93-9 (2)

83·2 (2) 89·6 (2) 90·2 (2)

94.4 (2)

107.5 (2)

113.8 (2)

Ge(5) - O(3)Ge(5) - O(4)

Ge(5)-O(5)

Ge(5) - O(7)Ge(5) - O(7A)

Ge(5)-O(8)

O(4)--Ge(1)--O(6) O(4)--Ge(1)--O(7) O(4)--Ge(1)--O(10)

O(4)—Ge(1)—O(11)O(4)—Ge(1)—O(8A)

O(6)-Ge(1)-O(7) O(6)-Ge(1)-O(10)

O(6) - Ge(1) - O(11)

O(6) - Ge(1) - O(8A)O(7) - Ge(1) - O(10)

O(7) - Ge(1) - O(11)

O(7)—Ge(1)—O(8A)O(10)—Ge(1)—O(11)

O(10)—Ge(1)—O(8A)O(11)—Ge(1)—O(8A)O(2)—Ge(2)—O(9)

O(2) - Ge(2) - O(10)

| O(2)-Ge(2)-O(12) | 113.5 (2) | O(1)—Li(2)—O(5B) | 173.0 (5) |
|-------------------------|-----------|---------------------------|-----------|
| O(9) - Ge(2) - O(10) | 110.7 (2) | O(2) - Li(2) - O(1B) | 99·9 (4) |
| O(9) - Ge(2) - O(12) | 103.8 (2) | O(2) - Li(2) - O(3A) | 87·3 (4) |
| O(10) - Ge(2) - O(12) | 107.2 (2) | O(2) - Li(2) - O(4B) | 178.8 (7) |
| O(1) - Ge(3) - O(2) | 86·1 (2) | O(2) - Li(2) - O(5B) | 103.1 (5) |
| O(1) - Ge(3) - O(3) | 92·1 (2) | O(1B)— $Li(2)$ — $O(3A)$ | 170.7 (5) |
| O(1) - Ge(3) - O(4) | 91·8 (2) | O(1B) - Li(2) - O(4B) | 79.5 (4) |
| O(1)-Ge(3)-O(6) | 173.8 (2) | O(1B) - Li(2) - O(5B) | 90.4 (5) |
| O(1) - Ge(3) - O(1A) | 83·6 (2) | O(3A) - Li(2) - O(4B) | 93.3 (4) |
| O(2)—Ge(3)—O(3) | 171.8 (2) | O(3A)—Li(2)— $O(5B)$ | 93.8 (4) |
| O(2) - Ge(3) - O(4) | 91.5 (2) | O(4B) - Li(2) - O(5B) | 77.9 (4) |
| O(2)-Ge(3)-O(6) | 90.9 (2) | O(10) - Li(3) - O(6A) | 99.1 (4) |
| O(2)—Ge(3)— $O(1A)$ | 92.1 (2) | O(10) - Li(3) - O(8A) | 78.6 (4) |
| O(3)-Ge(3)-O(4) | 80.5 (2) | O(10) - Li(3) - O(10A) | 85.1 (4) |
| O(3)—Ge(3)—O(6) | 90.2 (2) | O(10) - Li(3) - O(11A) | 177.7 (5) |
| O(3) - Ge(3) - O(1A) | 95.6 (2) | O(10) - Li(3) - O(11B) | 96.2 (4) |
| O(4)-Ge(3)-O(6) | 82.9 (2) | O(6A)— $Li(3)$ — $O(8A)$ | 93·2 (4) |
| O(4) - Ge(3) - O(1A) | 173.9 (2) | O(6A)-Li(3)-O(10A) | 170.2 (6) |
| O(6) - Ge(3) - O(1A) | 101.9 (2) | O(6A)— $Li(3)$ — $O(11A)$ | 78.9 (3) |
| O(11) - Ge(4) - O(5A) | 113.3 (2) | O(6A)-Li(3)-O(11B) | 95.7 (5) |
| O(11) - Ge(4) - O(9B) | 107.5 (2) | O(8A)-Li(3)-O(10A) | 96.4 (5) |
| O(11) - Ge(4) - O(12A) | 111.3 (2) | O(8A)—Li(3)— $O(11A)$ | 100.4 (4) |
| O(5A)—Ge(4)— $O(9B)$ | 113-4 (2) | O(8A)— $Li(3)$ — $O(11B)$ | 170-3 (5) |
| O(5A)—Ge(4)— $O(12A)$ | 106.8 (2) | O(10A) - Li(3) - O(11A) | 97.0 (4) |
| O(9B)—Ge(4)— $O(12A)$ | 104-3 (2) | O(10A)—Li(3)— $O(11B)$ | 74.9 (3) |
| O(3)-Ge(5)-O(4) | 83.7 (2) | O(11A) - Li(3) - O(11B) | 85.1 (4) |
| O(3)-Ge(5)-O(5) | 93.0 (2) | O(1)—Li(4)—O(3) | 77-2 (4) |
| O(3)—Ge(5)—O(7) | 94·2 (2) | O(1)—Li(4)—O(5) | 87.6 (5) |
| O(3)-Ge(5)-O(8) | 103.7 (2) | O(1)-Li(4)-O(2A) | 79.9 (4) |
| O(3)— $Ge(5)$ — $O(7A)$ | 171.4 (2) | O(1)—Li(4)—O(9A) | 165-4 (6) |
| O(4) - Ge(5) - O(5) | 88·3 (2) | O(1) - Li(4) - O(12B) | 91.3 (5) |
| O(4)-Ge(5)-O(7) | 83.7 (2) | O(3)—Li(4)—O(5) | 79·2 (4) |
| O(4)-Ge(5)-O(8) | 170.7 (2) | O(3)-Li(4)-O(2A) | 88.6 (4) |
| O(4) - Ge(5) - O(7A) | 89.3 (2) | O(3) - Li(4) - O(9A) | 89.5 (5) |
| O(5)-Ge(5)-O(7) | 168-6 (2) | O(3)-Li(4)-O(12B) | 165.6 (7) |
| O(5)-Ge(5)-O(8) | 96.7 (2) | O(5)—Li(4)—O(2A) | 164.2 (6) |
| O(5) - Ge(5) - O(7A) | 91.7 (2) | O(5)-Li(4)-O(9A) | 96·1 (5) |
| O(7)-Ge(5)-O(8) | 90·1 (2) | O(5) - Li(4) - O(12B) | 91.8 (4) |
| O(7) - Ge(5) - O(7A) | 80.0 (2) | O(2A) - Li(4) - O(9A) | 93.9 (5) |
| O(8)— $Ge(5)$ — $O(7A)$ | 82.9 (2) | O(2A) - Li(4) - O(12B) | 98.0 (4) |
| O(3)—Li(1)—O(6) | 80·6 (4) | O(9A)—Li(4)— $O(12B)$ | 102.7 (4) |
| O(3)Li(1)O(7) | 80.2 (3) | | |



Fig. 2. The $Li_4Ge_5O_{12}$ structure viewed along the *a* axis. Oxygen anions lie in quasi-close-packed planes parallel to the ab* plane of the crystal. A center of symmetry is located at the center of the unit cell which is taken as the origin for labeling the atoms. Atoms labeled Li4', Ge5', O5' and O8' have coordinates given in Table 3 minus 'b'. Ge4' has the coordinates of Table 3 minus ٬c'.

These linked pairs join with like pairs in adjacent cells to form chains that run along the *a* axis of the crystal. Taken together the octahedra have $\langle Ge - O \rangle = 1.901 (51) \text{ Å}$, 1.821 < Ge - O < 2.000 Å and $O - Ge - O \ge 80.0^{\circ}$. The octahedra are edge linked to one another and lie in zigzag bands perpendicular to **a** and **b***. Each of the eight Li per unit cell is surrounded by an irregular oxygen octahedron. Overall $\langle \text{Li} - O \rangle = 2.132 (91) \text{ Å}$, 1.988 < Li - O < 2.315 Å and $O - \text{Li} - O \ge 74.9^{\circ}$. Four of the eight lithium per unit cell lie in rectangular channels that lie along the *a* axis of the crystal.



Fig. 3. Packing arrangement of Ge–O polyhedra in $Li_4Ge_5O_{12}$ viewed along the crystal *a* axis. Shaded triangles show cornerlinked Ge–O tetrahedra that run in chains along the *a* axis. Shaded rectangles show edge-linked Ge–O octahedra that lie in bands perpendicular to *a* and *b**. Four of the eight Li per unit cell lie in the open rectangular channels that lie along **a**. Centers of symmetry are marked by (\oplus) and O(2) and O(10) are labeled.



Fig. 4. (a) Portion of the phase diagram (Murthy & Ip, 1964) near 90 wt % GeO₂ and 1213 K. (b) The type of peritectic transformation near 90 wt % GeO₂ and 1213 K that is consistent with our data.

Li₂O–GeO₂ phase diagram

Based on a number of crystal-growth experiments we see no evidence of the existence of the congruently melting phase ' $Li_6Ge_8O_{19}$ ' reported by Murthy & Ip (1964) and illustrated in Fig. 4(a). We have however grown large crystals of the compound Li₄Ge₅O₁₂ under quasi-equilibrium conditions and believe it to be an equilibrium phase in the Li₂O-GeO₂ system. In addition our data suggest that $Li_4Ge_5O_{12}$ may form by a peritectic reaction of the type illustrated in Fig. 4(b). This is supported by the observation that Li₄Ge₅O₁₂ crystals were grown from a melt of composition $Li_6Ge_8O_{19}$ and the observation that crystals of $Li_4Ge_5O_{12}$ could not be grown from a melt of composition Li₄Ge₅O₁₂. Furthermore, the reported eutectic temperatures in this region are uncertain to ± 10 K.

Differential thermal analysis (DTA) and roomtemperature X-ray diffraction experiments were performed to determine which phases exist in equilibrium at temperatures and compositions close to those used during crystal growth. The DTA measurements were abandoned because our minimum heating and cooling rates were too rapid for equilibrium to be attained. In addition, considerable thermal hysteresis and the presence of metastable side reactions complicate the thermograms. Problems similar to these were reported by West & Blake (1976) for the Li₂GeO₃-Na₂GeO₃ system. X-ray diffractograms of rapidly cooled samples indicated that they contained non-equilibrium phase mixtures typical of samples that are non-uniformly cooled. Additionally it should be noted that there is a strong tendency towards glass formation for compositions above 65 wt % GeO_2 (Shelby & Ruller, 1987).

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Crystal Chemistry of *cyclo*-Hexaphosphates. VIII. Structure of Hydroxylammonium *cyclo*-Hexaphosphate Tetrahydrate

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Abstract. (NH₃OH)₆P₆O₁₈.4H₂O, $M_r = 750 \cdot 11$, triclinic, $P\overline{1}$, $a = 10 \cdot 365 (5)$, $b = 9 \cdot 278 (4)$, $c = 7 \cdot 280 (3)$ Å, $\alpha = 108 \cdot 39 (5)$, $\beta = 100 \cdot 30 (5)$, $\gamma = 96 \cdot 02 (5)^\circ$, $V = 643 \cdot 8$ Å³, Z = 1, $D_x = 1 \cdot 934$ Mg m⁻³, λ (Mo $K\alpha$) = 0 $\cdot 7107$ Å, $\mu = 0 \cdot 553$ mm⁻¹, F(000) = 388, T = 298 K, final $R = 0 \cdot 028$ for 4750 reflections. The centrosymmetric P₆O₁₈ ring anion is located at the centre of the unit cell. Three crystallographically independent hydroxylammonium groups co-exist in the atomic arrangement. The three-dimensional hydrogen-bond network is described.

Introduction. Investigation of organic cationcyclo-hexaphosphates has been recently initiated by the characterization of the tris(ethylenediammonium) cyclo-hexaphosphate, $(NH_3-CH_2-CH_2-NH_3)_3$ -P₆O₁₈.2H₂O (Durif & Averbuch-Pouchot, 1989). In the present work we describe the preparation and crystal structure of a second example for such compounds, the hydroxylammonium cyclohexaphosphate tetrahydrate, $(NH_3OH)_6P_6O_{18}.4H_2O$.

Experimental. Crystals of the title compound have been prepared by using a metathesis reaction similar to that described by Boulle (1938) for the preparation of water soluble *cyclo*-triphosphates. Here the starting material is $Ag_6P_6O_{18}$.H₂O recently characterized by Averbuch-Pouchot (1989). Schematically the reaction is:

$$Ag_6P_6O_{18} + 6(OH - NH_3)Cl \rightarrow (OH - NH_3)_6P_6O_{18} + 6AgCl$$

Crystals obtained by slow evaporation of an aqueous solution at room temperature have various morphologies: stout triclinic prisms or thick diamond-like plates. The title compound is stable at

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room temperature. Crystal size: $0.40 \times 0.40 \times$ 0.35 mm. Density not measured. Philips PW1100 diffractometer, graphite monochromator. 19 reflections $(12.0 < \theta < 18.0^{\circ})$ for refining unit-cell dimensions. ω scan, scan width 1.20°, scan speed 0.02° s⁻¹. Total background measuring time 6 s. 5532 reflections collected $(3 < \theta < 35^\circ)$, $\pm h$, $\pm k$, l, $h_{\text{max}} = 16$, $\sin\theta/\lambda = 0.81 \text{ Å}^{-1}$. Two $k_{\rm max} = 14,$ $l_{\rm max} = 11.$ orientation and intensity control reflections (701 and $70\overline{1}$) measured every four hours without any significant variation. 5169 reflections obtained after averaging Friedel pairs ($R_{int} = 0.008$). Lorentz and polarization corrections, no absorption correction. Structure solved by direct methods (MULTAN78, Main, Hull, Lessinger, Germain, Declercq & Woolfson, 1978). H atoms located by difference Fourier syntheses. Anisotropic full-matrix least-squares refinement (on F), isotropic for H atoms. Unit weights. Final refinements with 4750 reflections corresponding to $I > 4\sigma_r$. Final R = 0.028 (wR = 0.033), S = 0.438, max. $\Delta/\sigma = 0.07$, max. peak height in the final difference Fourier synthesis = $0.425 \text{ e} \text{ Å}^{-3}$. Secondary extinction not refined. Scattering factors for neutral atoms and f', f'' from International Tables for X-ray Crystallography (1974, Vol. IV, Table 2.2B). Enraf-Nonius (1977) SDP used for all calculations, on a MicroVAX II computer.

Discussion. Table 1 reports the final atomic coordinates.*

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^{*} Lists of structure factors, anisotropic thermal parameters and H-atom parameters have been deposited with the British Library Document Supply Centre as Supplementary Publication No. SUP 53118 (42 pp.). Copies may be obtained through The Technical Editor, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.